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Passive scalar transport by a non-Gaussian turbulent flow (Batchelor regime)

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Passive scalar transport

(toy model for turbulence, anomalous scaling)

$$\frac{\partial \theta}{\partial t} + \mathbf{V} \cdot \nabla \theta = \kappa \Delta \theta$$

\mathbf{V} is a random velocity field

Kraichnan model: \mathbf{V} δ -correlated in time, Gaussian
and scale-invariant in space

- well studied, many results. But far from reality (no energy cascade)

$$B = \int A dt \xrightarrow{?} B_G \quad \begin{aligned} \langle B_G \rangle &= \langle B \rangle = \langle A \rangle t \\ D(B_G) &= D(B) = D(A) t \end{aligned}$$

$$\langle B^n \rangle \approx \langle B_G^n \rangle \quad \text{but} \quad \langle e^{nB} \rangle \neq \langle e^{nB_G} \rangle !$$

Non-Gaussian velocity field

Batchelor limit: $A_{ij} = \frac{\partial V_i}{\partial r_j}$ (viscous range, $\kappa \ll \nu$)

A_{ij} - stationary random process; Lagrangian frame

Balkovsky, Fouxon 1999: $\langle \theta^\alpha \rangle = e^{\gamma_\alpha t}$

Saturation of γ_α at large α .

No relation between θ and A_{ij} statistics.

Il'yn, Sirota, Zybin 2017 (submitted to PRE):

Exact expressions for γ_α in terms of A_{ij} statistics.

Universal saturation: $\gamma_\alpha = \text{const}$ if $\alpha \geq 2$

$$\text{(quasi)-Lagrangian frame: } \frac{\partial \theta}{\partial t} + \frac{\partial \theta}{\partial r_i} A_{ij} r_j = \varkappa \Delta \theta \quad \text{tr } A = 0$$

$A_{ij}(t)$ -stationary random process with known statistics

(for δ -correlated velocity, statistics of $A_{ij}(t)$ is the same in Euler and Lagrangian frames)

$$\begin{array}{ccc} \theta(\mathbf{r}, t) & \xrightarrow{\text{Fourier}} & \theta(\mathbf{k}, t) \\ \downarrow \mathbf{r}' = \mathbf{Q}\mathbf{r} & & \downarrow \mathbf{k} = \mathbf{p}\mathbf{Q} \\ \theta(\mathbf{r}', t) & \xrightarrow{\text{Fourier}} & \theta(\mathbf{p}, t) \end{array}$$

$$\dot{\mathbf{Q}} = -\mathbf{Q}\mathbf{A} ,$$

$$\mathbf{Q}(0) = \mathbf{I}$$

$$\frac{\partial}{\partial t} \theta(\mathbf{p}, t) = -\varkappa p_m (\mathbf{Q}\mathbf{Q}^T)_{mn} p_n \theta$$

$$\theta(\mathbf{p}, t) = \theta(\mathbf{p}, 0) e^{-\varkappa p_m \int (\mathbf{Q}\mathbf{Q}^T)_{mn}(t') dt' p_n}$$

$$\theta(t) \equiv \theta(\mathbf{r} = 0, t) = \int \theta(\mathbf{p}, t) d^3 p$$

$$\theta(\mathbf{p}, t) = \theta(\mathbf{p}, 0) e^{-\kappa \int p_m \left(\mathcal{Q} \mathcal{Q}^T \right)_{mn} (t') dt' p_n}$$

Isolated blob: $\theta(\mathbf{p}, 0) = e^{-l^2 \mathbf{p}^2}$

$$\theta(t) = \int \theta(\mathbf{p}, t) d^3 p = \int e^{-p_m D_{mn} p_n} d^3 p = (\det D)^{-1/2}$$

where

$$D_{mn} = \kappa \int \left(\mathcal{Q} \mathcal{Q}^T \right)_{mn} (t') dt' + l^2 \delta_{mn}$$

$$\langle \theta^\alpha \rangle = \langle (\det D)^{-\alpha/2} \rangle$$

Homogenous initial condition:

$$\langle \theta(\mathbf{p}, 0) \theta(\mathbf{p}', 0) \rangle = e^{-l^2 \mathbf{p}^2} \delta(\mathbf{p} - \mathbf{p}')$$

$$\langle \theta^\alpha \rangle = \langle (\det D)^{-\alpha/4} \rangle$$

Evolution matrix: Iwasawa decomposition

$$Q = Z \ d \ R \quad \begin{aligned} Z_{ij} &= 0 \quad \text{if } i > j, & Z_{jj} &= 1 \\ d &= \text{diag}(d_1, \dots, d_N) , \quad d_j > 0 , \quad d_1 \dots d_N &= 1 \\ R &\in SO(3) \end{aligned}$$

$$\dot{Q} = -QA , \quad Q(0) = I$$

Furstenberg 1963, Tutubalin 1978,...

$$t \rightarrow \infty \quad \begin{aligned} \exists \quad Z_\infty &= \lim_{t \rightarrow \infty} Z & Z_\infty \text{ depends on the realization} \\ \exists \quad \lambda_i &= \lim_{t \rightarrow \infty} \frac{\ln d_i(t)}{t} , \quad \lambda_1 \leq \lambda_2 \leq \dots \leq \lambda_N & \lambda_i \text{ are universal} \\ && \text{for given process} \end{aligned}$$

$$D_{mn} = \varkappa \int (QQ^T)_{mn}(t') dt' + l^2 \delta_{mn} = \varkappa \int (\mathbf{Z} \mathbf{d}^2 \mathbf{Z}^T)(t') dt' + l^2 I \quad (R \text{ vanishes!})$$

$$t \rightarrow \infty : \quad \det D = (\varkappa d^2 + C(l^2, \varkappa)) \propto d_1^{2\mu_1} \dots d_N^{2\mu_N} \quad \mu_j = \begin{cases} 1 & \text{if } d_j \gg 1 \\ 0 & \text{if } d_j \leq 1 \end{cases}$$

$$\langle \theta^\alpha \rangle = \langle (det D)^{-\alpha/2} \rangle \propto \left\langle d_1^{-\alpha \mu_1} \dots d_N^{-\alpha \mu_N} \right\rangle$$

Cumulant functions and functionals

Random variable $z \longrightarrow e^{W_z(k)} = \langle e^{ikz} \rangle = \int P(z) e^{ikz} dz$

Random process $\rho(t) \longrightarrow e^{W_\rho[\eta(t)]} = \left\langle e^{i \int \text{tr} \eta(t) \rho(t) dt} \right\rangle$

$$W[\eta(t)] = \int w(\eta(t)) dt \quad \text{for slow } \eta(t)$$

$$\frac{\ln d_i}{T} = z_i = \frac{1}{T} \int_0^T \rho_i dt \longrightarrow W_z(k) = w_\rho(k/T) \cdot T$$

$$z_i \xrightarrow{t \rightarrow \infty} \langle \rho_i \rangle = \lambda_i$$

$$\langle \theta^\alpha \rangle \propto \left\langle d_1^{-\alpha \mu_1} \dots d_N^{-\alpha \mu_N} \right\rangle \propto \left\langle e^{-\alpha (\mu_1 z_1 + \dots + \mu_N z_N) T} \right\rangle$$

$$\mu_j(z) = \begin{cases} 1 & \text{if } z_j > 0 \\ 0 & \text{if } z_j \leq 0 \end{cases}$$

$$\langle \theta^\alpha \rangle \propto \left\langle e^{-\alpha(\mu_1 z_1 + \dots + \mu_N z_N)T} \right\rangle \quad \mu_j(z) = \begin{cases} 1 & \text{if } z_j > 0 \\ 0 & \text{if } z_j \leq 0 \end{cases} \quad z_i = \frac{1}{T} \int_0^T \rho_i dt \xrightarrow{T \rightarrow \infty} \langle \rho_i \rangle = \lambda_i$$

Most naïve consideration: $z_i \approx \lambda_i$ $\lambda_1 \leq \lambda_2 \leq \dots \leq \lambda_N$

$$\langle \theta^\alpha \rangle \sim e^{-\alpha(\lambda_K + \dots + \lambda_N)T} \quad (\text{valid only for very small } \alpha)$$

Less naïve consideration: $\text{sign } z_i = \text{sign } \lambda_i$ $e^{W_\rho[\eta(t)]} = \left\langle e^{i \int \text{tr} \eta(t) \rho(t) dt} \right\rangle$

$$\langle \theta^\alpha \rangle \sim \left\langle e^{-\alpha(z_K + \dots + z_N)T} \right\rangle = e^{w_\rho(0, \dots, 0, i\alpha, \dots, i\alpha)T}$$

(valid only for α small enough)

Right consideration:

$$w_\rho(k) \xrightarrow{\text{saddle point}} P(z) \xrightarrow{\text{saddle point}} \langle \theta^\alpha \rangle \propto e^{\gamma_\alpha T},$$

$$\gamma_\alpha = \max_{k^*} \phi(k^*) , \quad \phi = -ik_j^* z_j + w_\rho(k^*) - \alpha \sum_{j=1}^N \mu_j(z) z_j , \quad z_j = -i \frac{\partial w_\rho}{\partial k_j} \Big|_{k^*}$$

$z_i \rightarrow \lambda_i$ but for some realizations $z_i(T) \leq 0$ even if $\lambda_i > 0$.

The bigger α , the more is their contribution to $\langle \theta^\alpha \rangle$.

$$\langle \theta^\alpha \rangle \propto e^{\gamma_\alpha T} \quad \mu_j(z) = \begin{cases} 1 & \text{if } z_j > 0 \\ 0 & \text{if } z_j \leq 0 \end{cases}$$

$$\gamma_\alpha = \max_{k^*} \phi(k^*) , \quad \phi = -ik_j^* z_j + w_\rho(k^*) - \alpha \sum_{j=1}^N \mu_j(z) z_j , \quad z_j = -i \frac{\partial w_\rho}{\partial k_j} \Big|_{k^*}$$

Inside the regions of constant μ_j :

$$\frac{\partial \phi}{\partial k_j^*} = 0 \quad \longrightarrow \quad k_j^* = i\alpha\mu_j \quad (= \text{'less na e consideration'})$$

$$\text{On the boundary } z_n = 0 : \quad \frac{\partial w_\rho}{\partial k_n} \Big|_{k^*} = 0 , \quad j \neq n : \quad k_j^* = i\alpha\mu_j$$

In both cases,

$$\langle \theta^\alpha \rangle = e^{w_\rho(k^*)T}$$

Relation between A and ρ

Il'yn, Sirota, Zybin 2016:

$$\dot{Q} = -QA \quad \longrightarrow \quad -A = Q^{-1}\dot{Q} = R^T d^{-1} Z^{-1} \frac{d}{dt}(ZdR)$$

Convenient variable: $X = RQ^{-1}\dot{Q}R^T = \underbrace{d^{-1}\dot{d}}_{\text{diagonal}} + d^{-1}Z^{-1}\dot{Z}d + \dot{R}R^T$

Then $X = -RAR^T$ $X_{nn} = -d_n^{-1}\dot{d}_n = \rho_n$

$$w_X(k) = w_A(ik_0 - k) - w_A(ik_0)$$

$$(k_0)_{ij} = \frac{2j-1-N}{2} \delta_{ij}$$

$$w_\rho(k_1, \dots, k_N) = w_X(k) \Bigg|_{k_{ij} = \begin{cases} k_i, & i = j \\ 0, & i \neq j \end{cases}}$$

Three-dimensional case

$$\xi = i(k_1^* - k_2^*) , \quad \eta = i(k_1^* - k_3^*) , \quad \psi = i(k_1^* + k_2^* + k_3^*)$$

Incompressibility $\longrightarrow \partial w_\rho / \partial \psi = 0 , \quad w_\rho = w_\rho(\xi, \eta)$

$$z_1 = \frac{\partial w_\rho}{\partial \xi} + \frac{\partial w_\rho}{\partial \eta}, \quad z_2 = -\frac{\partial w_\rho}{\partial \xi}, \quad z_3 = -\frac{\partial w_\rho}{\partial \eta}$$

$$z_j = -i \frac{\partial w_\rho}{\partial k_j^*}$$

Inside the regions of constant μ_j :

$$k_j^* = i\alpha\mu_j , \quad \xi^* = \alpha(\mu_2 - \mu_1) , \quad \eta^* = \alpha(\mu_3 - \mu_1)$$

$$\mu_j(z) = \begin{cases} 1 & \text{if } z_j > 0 \\ 0 & \text{if } z_j \leq 0 \end{cases}$$

On the boundary $z_2 = 0$: $\frac{\partial w_\rho}{\partial \xi^*} = 0 , \eta^* = \alpha(\mu_3 - \mu_1)$

$$w_\rho(\xi, \eta) = w_A(1-\xi, 2-\eta) - w_A(1, 2)$$

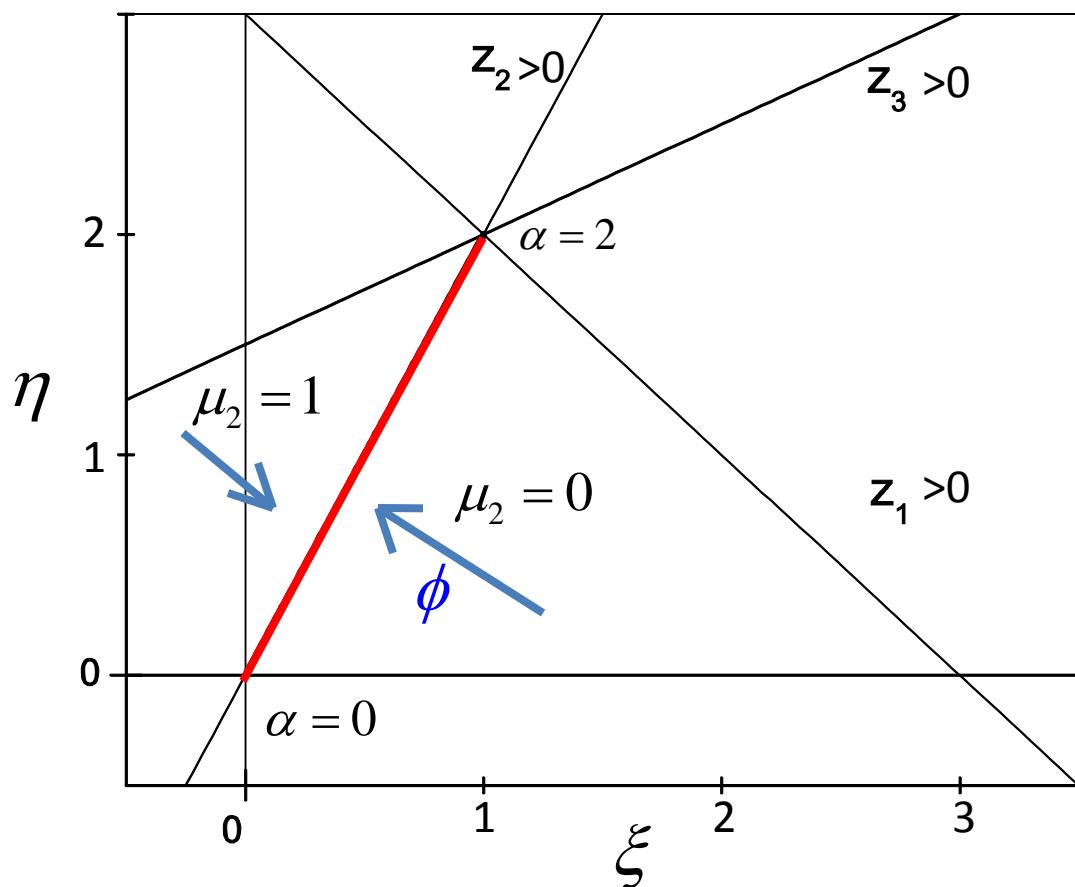
Examples: Gaussian statistics

$$w_A^G = D \left(\frac{1}{3} (k_1 + k_2 + k_3)^2 - (k_1^2 + k_2^2 + k_3^2) \right) = \frac{2}{3} D (\xi^2 + \eta^2 - \xi \eta)$$

$$w_\rho(\xi, \eta) = w_A(1-\xi, 2-\eta) - w_A(1, 2)$$

$$w_A(0,0) = w_\rho(0,0) = 0$$

$$\min_{\xi, \eta} w_\rho : \xi = 1, \eta = 2$$



$$\gamma_\alpha = \max_{\xi, \eta} \phi(\xi, \eta)$$

Vicinity of the boundary:

$$\mu_1 = 0, \quad \mu_3 = 1$$

$$\xi^* = \alpha \mu_2,$$

$$\eta^* = \alpha$$

$\alpha = 2$ -saturation

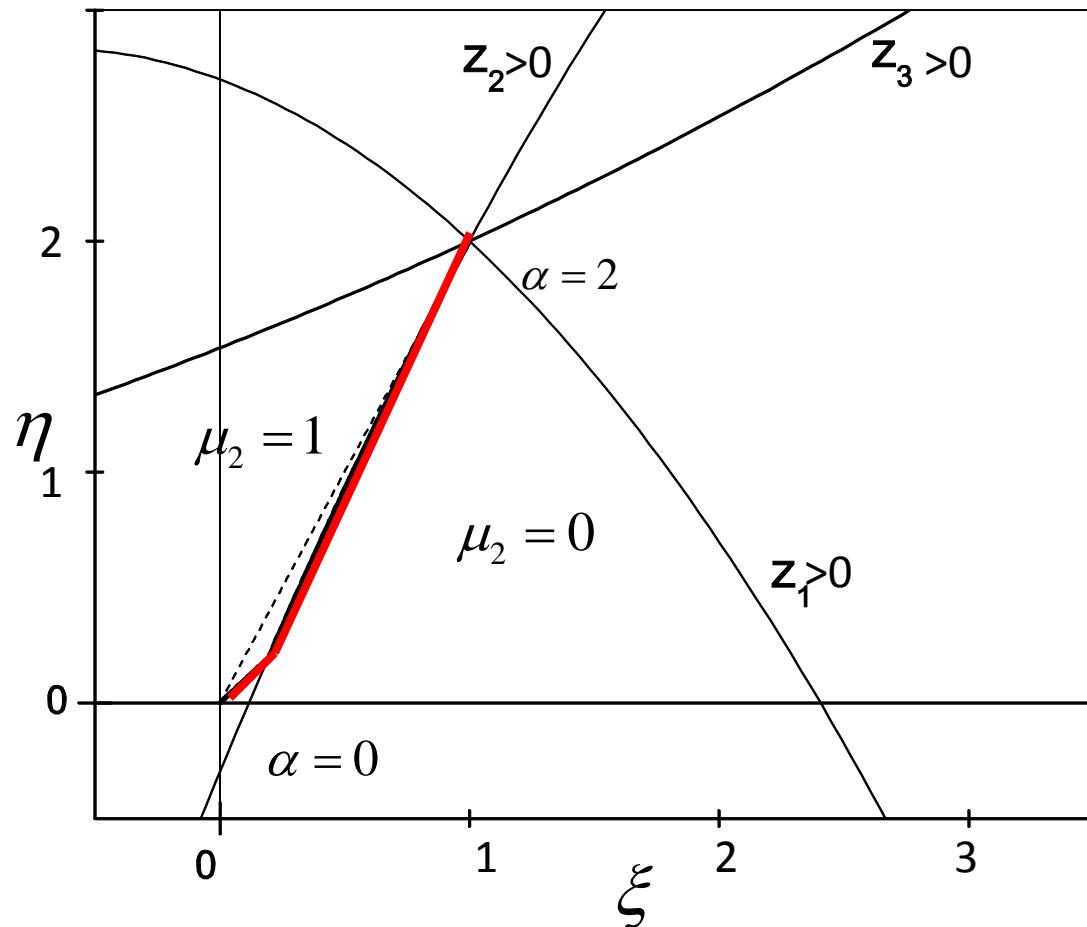
Examples: Small deviation from Gauss

$$w_A = \frac{2}{3}D(\xi^2 + \eta^2 - \xi\eta) - \frac{2}{9}F(\xi^3 + \eta^3) + \frac{1}{3}F\xi\eta(\xi + \eta), \quad F > 0$$

$$w_\rho(\xi, \eta) = w_A(1-\xi, 2-\eta) - w_A(1, 2)$$

(b)

$F > 0$



$$w_A(0,0) = w_\rho(0,0) = 0$$

$$\min w_\rho : \xi = 1, \eta = 2$$

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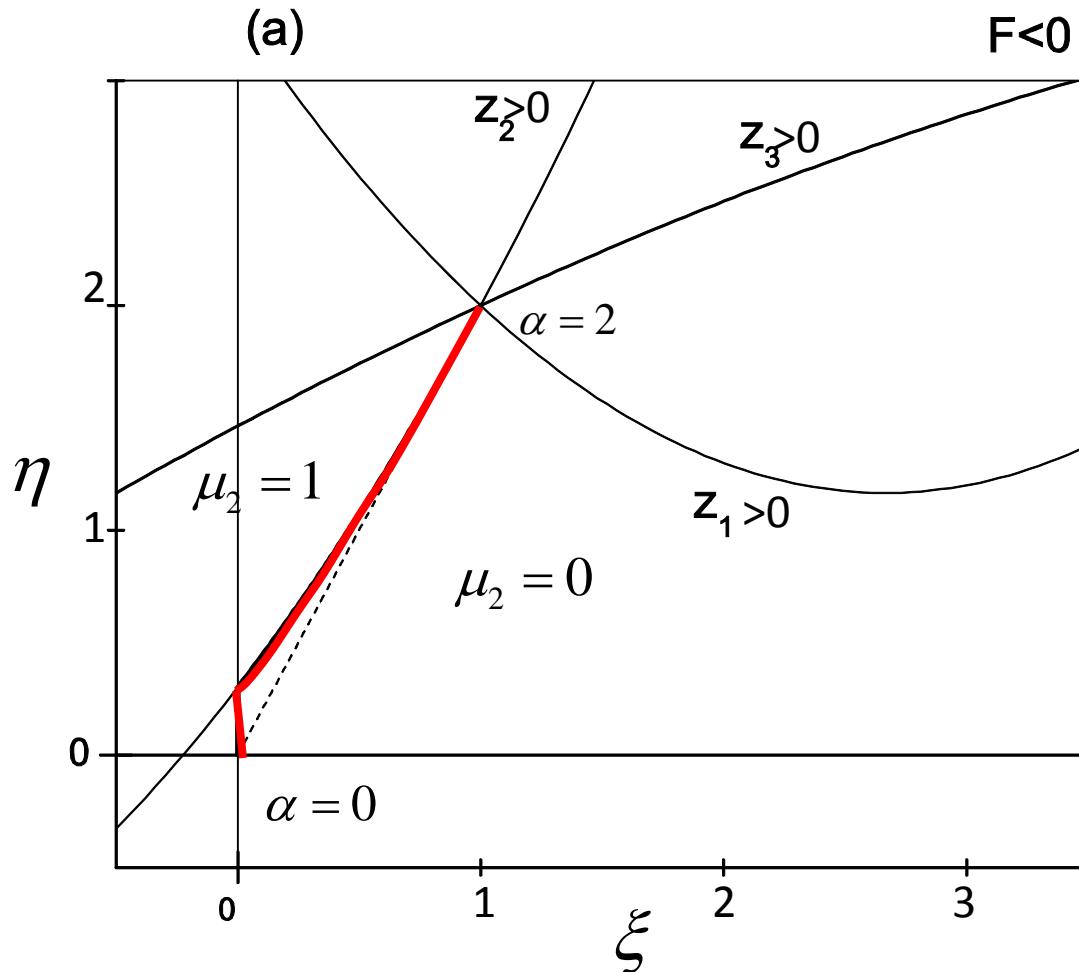
$$\gamma_\alpha = -2D + \frac{1}{2}D(\alpha - 2)^2 \\ - \frac{3}{32}D\left(\frac{F}{D}\right)^2 (\alpha - 2)^4 + O\left((F/D)^4\right)$$

$\alpha = 2$ -saturation
is universal!

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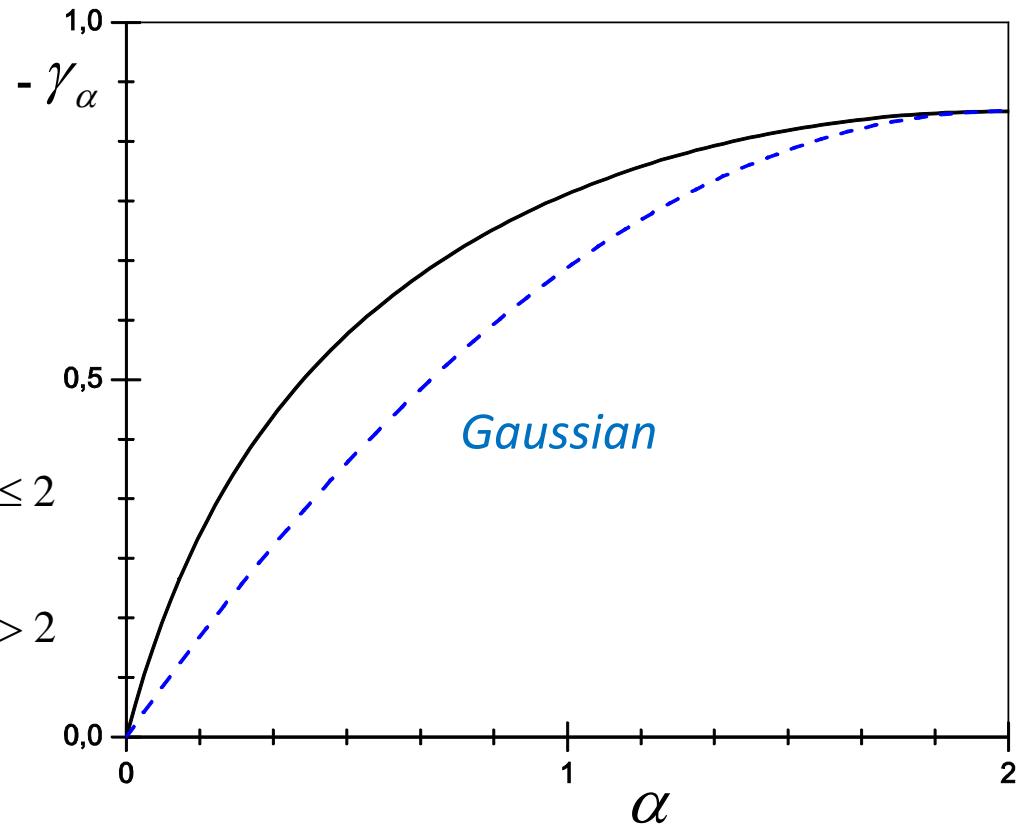
Examples: Exponential statistics

$$f(A_1, A_2, A_3) = f_0 e^{-c\sqrt{A_1^2 + A_2^2 + A_3^2}} \delta(A_1 + A_2 + A_3)$$

$$w_A = -\frac{1}{2} \ln \left(3c^2 - 4(\xi^2 + \eta^2 - \xi\eta) \right) + \frac{1}{2} \ln (3c^2)$$

$$w_\rho(\xi, \eta) = w_A(1 - \xi, 2 - \eta) - w_A(1, 2)$$

$$\gamma_\alpha = \begin{cases} \phi(\alpha/2, \alpha) = \frac{1}{2} \ln \frac{c^2 - 4}{c^2 - (\alpha - 2)^2}, & \alpha \leq 2 \\ \phi(1, 2) = \frac{1}{2} \ln \frac{c^2 - 4}{c^2}, & \alpha > 2 \end{cases}$$



Summary

- We analyze the passive scalar advection in a turbulent flow at time much bigger than the correlation time of the flow, and scales much smaller than the viscous scale (viscous range); diffusivity is assumed to be much smaller than viscosity, linear approximation for the velocity field is valid (Batchelor regime).
- We obtain exact expressions for the exponents γ_α of the Lagrangian scalar density moments in terms of Lagrangian velocity strain tensor statistics. (If velocity is assumed to be δ -correlated, the latter coincides with the Eulerian strain tensor statistics).
- The exponents saturate at the universal value $\alpha=2$, independently of the statistics. In the range between $\alpha=2$ and $\alpha=2$ they can differ significantly from those in *the Gaussian case*.